for C<sub>H</sub> strain

$$\frac{a_{1}}{a_{2}} = a \int \frac{1}{3} (1,0,0)$$

$$\frac{a_{2}}{a_{2}} = a \int \frac{1}{3} (-1/2, \frac{1}{3}/2, 0)$$

$$\frac{c}{a_{2}} = c \int \frac{1}{3} (0,0, \frac{1}{3} - 1)$$
and  $C_{H} = \left(\frac{d^{2}W}{d \frac{1}{3} - 2}\right) \int \frac{(9/2)}{\frac{1}{3} - 1}$ 

A further condition is obtained by considering the first derivatives of the energy with respect to the strains. The conditions for equilibrium with no applied stress are

$$\left(\frac{dW}{d\eta}\right)_{\eta=1} = 0$$
 and  $\left(\frac{dW}{dg}\right)_{g=1} = 0$ 

In the case of the  $C_{66}$  strain all three contributions have first derivatives which are independently zero; in the case of the  $C_{\rm H}$  strain the Coulomb and the full zone Fermi first derivatives yield negative contributions and must be matched by a positive contribution from the overlap-hole term to comply with the equilibrium condition, which thus becomes an independent and useful condition.

A  $C_{l_ll_l}$  type strain lowers the symmetry of the crystal such that the calculations become quite involved. No attempt was made to account for the  $C_{l_ll_l}$  shear constant.